# The Crystal Structure of the a Form of Anhydrous Copper 8-Hydroxyquinolinate

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The crystal structure of the  $\alpha$  form of anhydrous copper 8-hydroxyquinolinate, Cu(C<sub>9</sub>H<sub>6</sub>NO)<sub>2</sub>, has been determined by the Weissenberg equi-inclination X-ray method. The crystals, of bronze colored, fibrous morphology, were found to be monoclinic in the space group  $P2_1/c$ . The unit cell of dimensions  $a=7.40\pm0.03$ ,  $b=3.84\pm0.02$ ,  $c=24.37\pm0.09$  Å, and  $\beta=98^{\circ}6'\pm10'$ , contains two molecules. Intensities were measured visually, and the data refined by difference synthesis and least squares with a final R index of 12.2% including all observed reflections. Absorption corrections were neglected.

The copper atoms, which lie on symmetry centers, are six-coordinated. They are bound to the two hydroxyquinolinate anions through the ring nitrogen atoms and the hydroxyl oxygen atoms in transplanar configurations, and to the oxygen atoms from adjacent molecules; they reside at the centers of distorted octahedra. The anion rings are essentially planar with the planes of the opposing anions parallel but slightly displaced from each other. The copper and oxygen atoms are displaced from both of these planes.

## Introduction

Among known metal ion chelating agents is 8-hydroxyquinoline, which forms bidentate chelates with most of the divalent metals. Although these salts are well known, the nature of their anion-cation bonds is not well understood. Charles, Freiser, Fridel, Hillard & Johnson (1956) studied the infrared spectra of these complexes and found that they fell into two distinct groups when the positions of their carbon-oxygen stretching vibrations  $(9\mu)$  were plotted against the atomic weights of the metals. One group, the transition metals, fell on a line of very steep positive slope while the other, the non-transition metals, resided on an almost horizontal line of slight negative slope. It was postulated that this result was due either to a difference in structure between the two groups, or to involvement of the 3d orbitals in the metal-ligand bonding for the transition metal chelates. This question, however, was not resolved.

The cupric salt of 8-hydroxyquinoline,  $Cu(C_9H_6NO)_2$ , has been shown to be dimorphic in both the anhydrous and dihydrate forms. Suito & Sekido (1957) described the  $\alpha$  and  $\beta$  forms of the dihydrate. The  $\alpha$  form, of fibrous morphology, was metastable and transformed into the  $\beta$  form which was plate-like, of tetragonal, hexagonal, or rhombohedral shape. X-ray powder diffraction patterns of the two forms indicated a probable difference in structure since the  $\beta$  form displayed many more lines. Merritt, Cady & Mundy (1954) investigated the  $\beta$  form and found it to be isomorphous with zinc hydroxyquinolinate dihydrate. Jonassen (1962) has shown the anhydrous form to exist as a metastable  $\alpha$ form and a stable  $\beta$  form also. The anhydrous forms also exhibited a difference in their X-ray powder patterns and, in addition, an important difference in their infrared absorption spectra. When the positions of their  $9\mu$  absorption peaks were plotted against atomic weight, it was found that the  $\alpha$  form fell into one of the groups described by Charles *et al.*, the transition metal group, whereas the  $\beta$  form fell into the other. A complete structure determination of the anhydrous  $\beta$  form has been accomplished by Palenik (1964). Further analysis of these complexes depends on the structure determinations of the remaining members of this family of compounds. In this regard the structure determination of the  $\alpha$  form of the copper salt was undertaken.

### Experimental

The  $\alpha$  form of anhydrous copper 8-hydroxyquinolinate crystallizes in either lath-like or acicular crystals. Both types are essentially transparent with a yellowish gold tint, display well developed faces, and cleave normal to the needle axis. In addition transmitted light is polarized parallel to the needle axis.

Examination of single crystals of the compound with a Weissenberg camera using equi-inclination settings and copper radiation, revealed the following systematic absences:

hkl no absences,
h0l absent when l is odd,
0k0 absent when k is odd.

These absences are characteristics only of the space group  $P2_1/c$  ( $C_{2h}^5$ ). In addition, the unit-cell dimensions were found to be:

$$a = 7.40 \pm 0.03, \quad b = 3.84 \pm 0.02, \quad c = 24.37 \pm 0.09 \text{ Å}$$
  
 $\beta = 98^{\circ}6' \pm 10'.$ 

Since the density, while not measured precisely, was found to be slightly greater than  $1.6 \text{ g.cm}^{-3}$ , it was concluded that there were two molecules per unit cell, as calculation based on this arrangement yields a density of  $1.71 \text{ g.cm}^{-3}$ .

In order to avoid the necessity of correcting for absorption, two very small crystals were used in making the photographs for the intensity measurements. An acicular crystal approximately  $0.050 \times 0.058$  mm in cross section was used for *b*-axis photographs, and a fragment about  $0.03 \times 0.03$  mm was cut from a larger crystal for rotation around the *a* axis. The crystals were mounted on a Weissenberg camera and their diffraction intensities recorded with exposures ranging up to 100 hours. Five films were exposed simultaneously in the multiple film technique. For Cu K $\alpha$  radiation the transmission factor of the Kodak No-Screen film employed was approximately 23.3%. Intensities were measured visually by comparison with a standard scale prepared from the larger crystal.

Out of 2463 non-symmetry-related reflections accessible to film observation, 1144 were observed. All intensities were measured a minimum of three times and their average value taken as the best value. Approximate scale factors between layers were determined by comparing reflections common to layers of both axes. High order reflections were corrected for the  $\alpha_1 - \alpha_2$  split by multiplying partially resolved intensities by 1.2 and completely resolved intensities by 1.5 in a manner similar to that of Ghose & Fehlmann (1965). An additional 1319 reflections which were theoretically recorded on the film, but were too weak to be observed, were considered to be less than the minimum observable value at their respective angles. Contributions from these reflections were included in the final R value, but were excluded from the refinement process. Corrections for Lorentz and polarization factors were accomplished in the standard fashion.

The fact that there was a very short axis in the unit cell implied a planar molecule, generally lying in the x-z plane, and that the greatest information could be gained from projections parallel to this axis. Therefore, a two-dimensional Patterson projection map was constructed from the hol data. Although examination of this Patterson projection suggested sites for all the heavy atoms, a more cautious approach was undertaken. A series of structure factor calculations of just the h0l reflections was embarked upon in which atoms, starting with copper only, and with estimated isotropic temperature factors, were added to the summations one at a time to observe their effect on the residual. Absolute structure factors were approximately obtained by multiplying the observed structure factors,  $F_o$ , by a scale factor such that the sum of the absolute values of  $F_o$  equalled the sum of the absolute values of the calculated structure factors,  $F_c$ . After each addition better agreement was achieved. The following atomic scattering factor curves were used in these calculations: Cu<sup>2+</sup> from Watson & Freeman (1961), corrected for the real part of anomalous dispersion; oxygen and nitrogen from Hoerni & Ibers (1954); and carbon from McWeeny (1951). In addition, the isotropic temperature factor coefficients were assumed to be: 2.0 for copper and oxygen, 3.4 for nitrogen, and 4.0 for carbon.

After all but the hydrogen atoms were included, the method of successive Fourier synthesis was utilized to refine the x and z atom positions until no further sign changes were evident.

The y atom positions were determined by assuming that the atoms were in the x-z plane, and then estimating from a subsequent electron density projection around the [100] axis made from the 0kl reflections. Again the successive Fourier synthesis technique was used and more accurate y positions ascertained. At this time the gross features of the structure of the molecule had been determined, and the R value was about 30% for all the observed reflections. A careful check of all observations revealed several errors which were corrected. In addition the 002 reflection appeared to be a victim of secondary extinction, and therefore was omitted from later refinement and residual calculations.

At this time a least-squares refinement program by Busing, Martin & Levy (1962) written for the IBM 7090 computer was adapted to the IBM 7044. Two cycles of this program on all the observed data assuming unit weighting factors, which were used throughout refinement, reduced R to about 18%. The function chosen to be minimized by this program was  $\Sigma |F_o - sF_c| / \Sigma |F_o|$  where s is a scale factor. From the atom positions derived from this process, a difference map about the [010] axis was calculated and plotted. Estimates of the possible positions of the six hydrogen atoms in the asymmetric unit were made from this projection. Their y positions were assumed to be the same as that of the carbon atom to which they were bound. This projection also served to call attention to the need for anisotropic temperature factors; therefore, anisotropic temperature coefficients were incorporated of the form:

 $\exp\left\{-(h^2\beta_{11}+k^2\beta_{22}+l^2\beta_{33}+2hk\beta_{12}+2hl\beta_{13}+2kl\beta_{23})\right\}.$ 

Five more least-squares cycles were calculated and R fell to 12.3%, including contributions due to the hydrogen atoms whose positions were held constant. An isotropic temperature coefficient of 5.0 was assumed for the hydrogen atoms, and converted to anisotropic coefficients for the calculations. Difference maps about both axes were made from these positions. omitting the hydrogen atoms, so that new hydrogen positions could be determined. Probable positions were found for all six of the hydrogen atoms. With these new hydrogen positions, R was lowered to 12.2%, at which point refinement was terminated. However, when contributions from the unobserved reflections were added, R rose to  $15 \cdot 1\%$ . A list of all the structure amplitudes together with their final calculated counterparts is given in Table 1.

The resulting atomic coordinates of the heavier atoms are given in Table 2 in fractions of the cell edge, and their temperature parameters in Table 3, along with their estimated uncertainties, which do not include the uncertainties in the lattice constants. Table 1. Observed and calculated structure factors

Each group of three columns contains l,  $10F_o$ , and  $10F_c$  respectively, and is headed by values of h and k common to the group. Minus signs in the  $F_o$  columns denote unobserved reflections.

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Table 1 (cont.)

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Table 2. Positional parameters	of the heavy atoms	in fractional coordinates	and their estimated	standard deviations
-		<b>y</b>		

	x/a	y/b	z/c	$\sigma_x$	$\sigma_y$	$\sigma_z$
Cu	0.0000	0.0000	0.0000			
0	-0.0876	0.2523	0.0599	0.0011	0.0025	0.0003
Ν	0.2284	-0.0592	0.0524	0.0012	0.0027	0.0003
C(1)	0.3858	-0.5060	0.0443	0.0017	0.0032	0.0005
C(2)	0.5349	-0.2137	0.0877	0.0016	0.0034	0.0005
C(3)	0.5105	-0.1063	0.1390	0.0017	0.0032	0.0005
C(4)	0.3433	0.0200	0.1494	0.0015	0.0030	0.0004
C(5)	0.3043	0.1760	0.2005	0.0018	0.0032	0.0005
C(6)	0.1417	0.3288	0.2045	0.0017	0.0033	0.0005
C(7)	-0.0008	0.3567	0.1574	0.0018	0.0033	0.0005
C(8)	0.0311	0.2296	0.1058	0.0015	0.0030	0.0005
C(9)	0.2031	0.0733	0.1030	0.0012	0.0029	0.0004

Table 3. Least-squares thermal parameters and their estimated standard deviations

	$\beta_{11}$	$\beta_{22}$	$\beta_{33}$	$\beta_{12}$	$\beta_{13}$	$\beta_{23}$
Cu	0·0148	0·1012	0·0017	0·0066	0·0002	-0.0033
	(0·0005)	(0·0027)	(0·0001)	(0·0013)	(0·0001)	(0.0004)
0	0·0162	0·1295	0·0015	0·0019	-0.0001	-0.0040
	(0·0018)	(0·0101)	(0·0002)	(0·0035)	(0.0005)	(0.0011)
N	0·0181	0·0891	0·0012	-0.0012	0·0010	-0.0017
	(0·0021)	(0·0104)	(0·0002)	(0.0042)	(0·0005)	(0.0011)
<b>C</b> (1)	0·0193	0·0686	0·0017	0·0092	0·0003	-0.0007
	(0·0029)	(0·0114)	(0·0003)	(0·0048)	(0·0007)	(0.0014)
C(2)	0·0155	0·0845	0·0020	0·0072	0·0014	-0.0028
	(0·0027)	(0·0122)	(0·0003)	(0·0047)	(0·0007)	(0.0015)
C(3)	0·0205	0·0705	0·0015	0·0037	-0.0004	0.0011
	(0·0030)	(0·0127)	(0·0002)	(0·0046)	(0.0007)	(0.0013)
C(4)	0·0216	0·0548	0·0013	-0.0008	0·0013	0·0003
	(0·0025)	(0·0108)	(0·0002)	(0.0049)	(0·0006)	(0·0013)
C(5)	0·0228	0·0775	0·0014	0·0010	0·0006	-0.0002
	(0·0032)	(0·0011)	(0·0003)	(0·0047)	(0·0007)	(0.0013)
C(6)	0·0170	0·0831	0·0017	-0.0007	0·0010	-0.0006
	(0·0029)	(0·0117)	(0·0003)	(0.0044)	(0·0007)	(0.0014)
C(7)	0·0225	0·0843	0·0015	-0.0001	0·0024	-0.0013
	(0·0031)	(0·0122)	(0·0003)	(0.0045)	(0·0007)	(0.0013)
C(8)	0·0171	0·0565	0·0018	0·0093	0·0009	-0.0014
	(0·0026)	(0·0106)	(0·0003)	(0·0046)	(0·0006)	(0.0014)
C(9)	0·0164	0·0519	0·0015	0·0062	0·0004	-0.0010
	(0·0024)	(0·0110)	(0·0002)	(0·0042)	(0·0006)	(0.0011)

 Table 4. Probable hydrogen atom positions in fractional coordinates and distance from the associated carbon atom (in brackets)

	x/a	y/b	z/c	Distance
H(1)[C(1)]	0.357	0.246	0.000	1·08 Å
H(2)[C(2)]	0.615	-0.367	0.068	1.00
H(3)[C(3)]	0.613	-0.030	0.175	1.12
H(5)[C(5)]	0.407	0.198	0.236	1.07
H(6)[C(6)]	0.066	0.349	0.241	1.12
H(7)[C(7)]	-0.101	0.201	0.163	0.95

Probable hydrogen positions designated by the number of the carbon atoms to which they are attached, and their calculated bond lengths to the associated carbons are given in Table 4. From the atomic positions in Table 2, all the intramolecular distances and angles of interest were computed, and are given in Fig.1. Their maximum probable error in bond length was 0.018 Å, and  $1.2^{\circ}$  in bond angle. Their average probable error in bond length was 0.015 Å and in bond angle  $1.0^{\circ}$ .

# Discussion

Fig. 1 illustrates the final calculated bond lengths and angles for the molecule. Since there were no unusual intermolecular or intramolecular distances involved, the resulting structure is stereochemically sound and may be regarded as a molecular structure with only weak bonding, as signified by the relatively long bond length of 3.324 Å between the copper and oxygen atoms of adjacent molecules. Thus, the copper atoms in this crystal, unlike those of the  $\beta$  form investigated by Palenik (1964) which were fivefold coordinated in distorted tetragonal pyramidal forms, are sixfold coordinated. The existence of four short bonds and two long bonds in the distorted octahedral arrangement around the copper make this compound one more example of Jahn–Teller distortion in a Cu<sup>II</sup> complex. The two hydroxyquinolate anions that are bound through the hydroxyl oxygen and ring nitrogen atoms essentially form two planes parallel to each other, but displaced by a small distance. The best least-squares planes through various atoms of the asymmetric unit have been calculated following Scho-



Fig.1. Intermolecular distances and angles.

maker, Waser, Marsh & Bergman (1959) and are given in Table 5. Both the copper and oxygen atoms are slightly displaced from the planes of the quinoline rings.

The copper-oxygen distance of 1.939 Å found in this study is in agreement with that of the  $\beta$  form and those found in other studies. This distance is slightly shorter than the copper-nitrogen bond which was 1.984 Å and was discussed by Palenik in a similar situation for the  $\beta$  form. He pointed out that this difference, in this case 0.045 Å, can be accounted for by the Schomaker and Stevenson correction to the additivity of single-bond radii, and that the contention of Lingafelter, Simmons, Morosin, Scheringer & Freiburg (1961), that this inequality ought to be attributed to steric hinderance, should be re-examined. Although the agreement with the Schomaker and Stevenson correction is better for the  $\alpha$  form than the  $\beta$  form, the molecule also appears to be sensitive to steric repulsion and no conclusions can be reached in regard to this question.

A comparison of the ring structures of the  $\alpha$  and  $\beta$  forms reveals that they are essentially the same; the largest difference in the chemically equivalent bond is 0.023 Å with an average difference of 0.012 Å. This average difference is of the same order as the estimated standard deviations for the bond, and it is felt that this represents a good estimate of the accuracy of the structure.

In light of the difference in structure between the  $\alpha$ and  $\beta$  forms of anhydrous copper 8-hydroxyquinolinate, it is very evident that this is the cause of the difference between their infrared spectra. Since the  $\beta$  form (Palenik, 1964) and the zinc complex (Merritt, 1953) have both been shown to have a tetrahedral structure, and are in one of the groups studied by Charles *et al.* (1956), whereas the  $\alpha$  form is in the other, it seems reasonable to predict that future studies will show that the reason for the differences in the two groups of divalent metal chelates is a difference in structure.

	Table 5. Least-s	quares plan	es	
	Plane and devia	tions ( $\times 10^3$ )	from the	plane
~	т	ŤŤ Í	TTT	-

		Atom	1	11	111	
		Cu	- 093	-156	-112	
		0	095	049	090	
		Ν	-020	-052	-032	
		C(1)	-017	-045	-035	
		C(2)	071	068	060	
		C(3)	005	018	003	
		C(4)	-009	-001	-005	
		C(5)	-051	-027	-038	
		C(6)	-019	-002	-001	
		C(7)	007	-000	018	
		C(8)	029	005	030	
		C(9)	002	-014	-001	
	Direc	tion cosines	(×104)			
	1	with respect	to			
Plane	а	b	c*	Distance from	ı origin	Description of plane
I	3836	8919	-2397	0.0933	U	Asymmetric unit
II	3909	8927	- 2242	0.1562		Hydroxyquinoline anion
III	3821	8943	2329	0.1117		Quinoline ring only



Fig. 2. Views of the structure along [100] and [001].

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# The Crystal Structure of 10,10'-Dianthronyl

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The crystal structure of 10,10'-dianthronyl has been determined from three-dimensional X-ray data. The atomic coordinates and anisotropic temperature factors have been refined by the method of least squares. The dimer has twofold symmetry perpendicular to the C(10)-C(10') bond and the planes of the two half-molecules make an angle of 145°18' with each other. The thermal vibrations indicate libration of the two half-molecules about C(10)-C(10').

## Introduction

In the course of studies of the diamagnetic anisotropy of organic crystals Lasheen (1963, 1964) crystallized anthrone from acetic acid and from ethyl alcohol, and in both cases he obtained two types of crystal, which he assumed to be two modifications of anthrone. Srivastava (1957a,b, 1961) had determined a crystal